Week 3 (Lecture 2)

Computational Methods

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Overview

We already know the nature of PDE's, now we will attempt to discretize and compute the PDE's.

We will focus on model problems.

These model problems have difficulties that are shared with more realistic models, but much simpler to handle.

Overview (cont'd)

The model problems that will be discussed are

-
$$u_t + a u_x = 0$$
 hyperbolic

-
$$u_t = \kappa u_{xx}$$

parabolic

-
$$u_{xx} + u_{yy} = 0$$
 elliptic

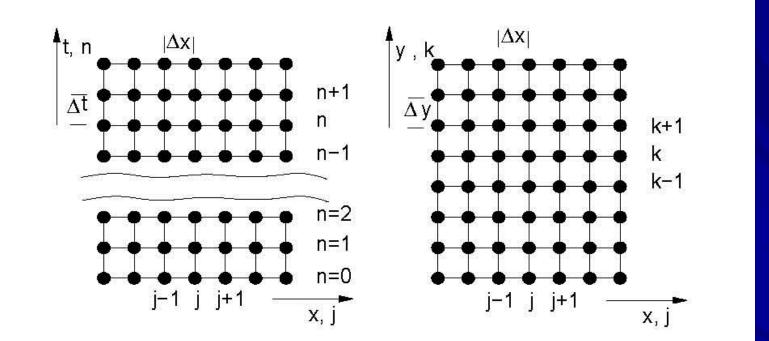
Overview (cont'd)

We will discretize the models using FD method.

Only simple uniform grids are used.

Stick to the notations introduced before.

Overview (cont'd)

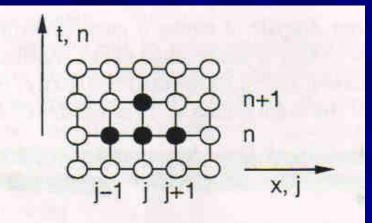


Evolution

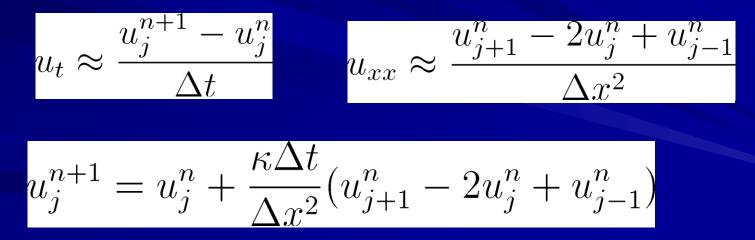
Equilibrium

Computation of Parabolic Equation

$$u_t = \kappa u_{xx}$$



Apply FTCS scheme



von Neumann (VN) Analysis

A method to determine stability of numerical schemes

Decompose solution in terms of Fourier modes

$$u_j^n = \operatorname{Real}(g^n \exp(ij\theta))$$

Variation in space in terms of sine wave from $\theta = [0, \pi]$

g is amplification factor, indication for stability

Even though it only applies for linear problems, it provides a reasonably accurate guide to more general cases.

Von Neumann Analysis on FTCS

Rewrite FTCS scheme solving 1D heat equation

$$u_j^{n+1} = u_j^n + \mu(u_{j+1}^n - 2u_j^n + u_{j-1}^n)$$

Using von Neumann analysis

$$g^{n+1}\exp(ij\theta) = g^n\exp(ij\theta) + \mu g^n(\exp(i(j+1)\theta) - 2\exp(i(j)\theta) + \exp(i(j-1)\theta))$$

$$g = 1 + \mu(\exp(i\theta) - 2 + \exp(-i\theta))$$

= $1 + \mu(\cos\theta + i\sin\theta - 2 + \cos\theta - i\sin\theta)$
= $1 + 2\mu(\cos\theta - 1)$

VN on FTCS (cont'd)

This means that if input is pure sine wave (Fourier mode), after 1 time step the sine wave will be amplified by g which depends on θ and μ.

Restriction must be done on \mu, not on \theta = [0, \pi].

$$-1 \le \cos\theta \le 1$$

For stability requires



$$u \le \frac{1}{2}$$

Comments using FTCS on 1D Heat Eqn

The scheme is conditionally stable

If we use small Δx , we need extremely small Δt since it is scaled in Δx^2

It can be shown that the restriction

$$u \le \frac{1}{2}$$

applies to all numerical schemes solving the 1D heat eqn

FTCS Solution on 1D Heat Eqn

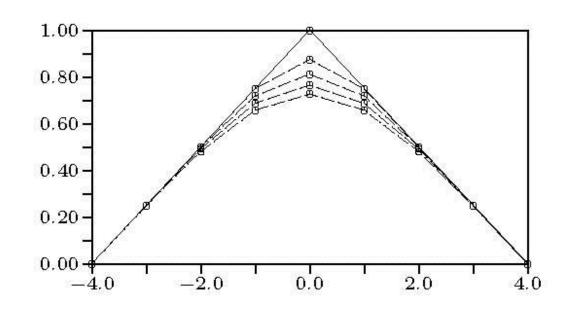
The electric blanket problem-heat added in the center of blanket

Use 1D model where at t=0, add heat such that

$$u(x) = (1+x), \text{if}(-1 \le x \le 0)$$
$$u(x) = (1-x), \text{if}(0 \le x \le 1)$$

And switch off power

FTCS Solution on 1D Heat Eqn (cont'd)

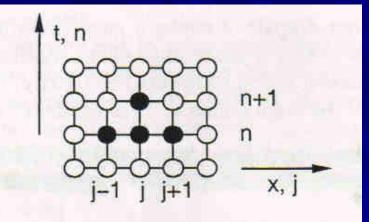


Results are very accurate but oscillations will grow wild if restriction is violated.

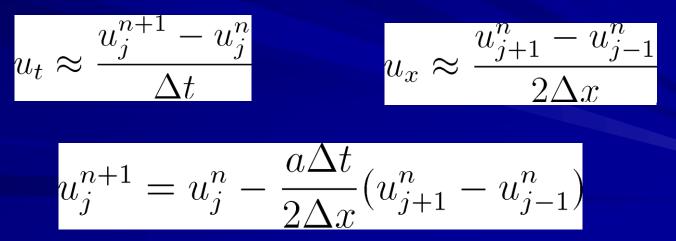
$$\mu > \frac{1}{2}$$

Computation of Hyperbolic Equation

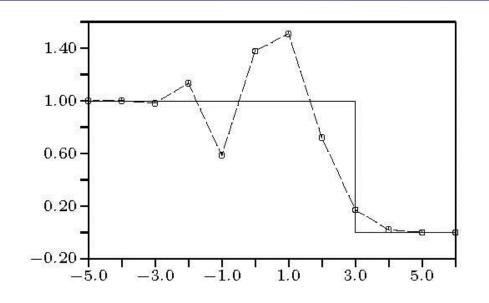
$$u_t + au_x = 0$$



Apply FTCS scheme



FTCS Solution on 1D Advection Eqn



FTCS is unstable

Just because it works for 1 type of PDE, does not mean it will work for other types

VN Analysis on FTCS (advection)

Rewrite FTCS scheme solving 1D heat equation

$$u_j^{n+1} = u_j^n + 0.5\nu(u_{j+1}^n - u_{j-1}^n)$$

Using von Neumann analysis

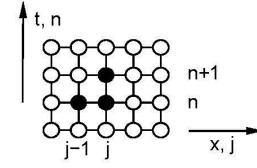
$$g^{n+1}\exp(ij\theta) = g^n \exp(ij\theta) + 0.5\mu g^n (\exp(i(j+1)\theta) - \exp(i(j-1)\theta))$$

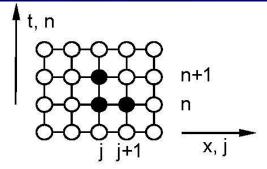
$$g = 1 - 0.5\nu(\exp(i\theta) - \exp(-i\theta))$$
$$= 1 - 0.5i\nu\sin\theta$$
$$|g| = (1 + \nu^2\sin^2\theta)^{0.5}$$

FTCS unconditionally unstable!

Computation of Hyperbolic Equation

$$u_t + au_x = 0$$

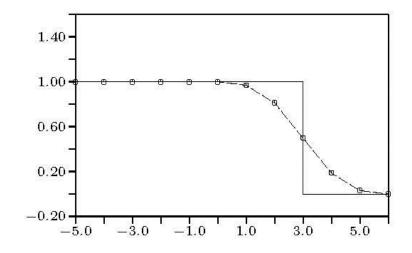




Apply 1st order upwind scheme

$$\begin{split} u_t &\approx \frac{u_j^{n+1} - u_j^n}{\Delta t} \\ u_x &\approx \frac{u_j^n - u_{j-1}^n}{\Delta x}, \text{if}(a > 0) \\ u_x &\approx \frac{u_{j+1}^n - u_j^n}{\Delta x}, \text{if}(a < 0) \\ u_x &\approx \frac{u_{j+1}^n - u_j^n}{\Delta x}, \text{if}(a < 0) \\ \end{split}$$

1st order Upwind Solution on 1D Advection Eqn



1st order upwind is conditionally stable

It can be shown using VN analysis that scheme is stable if

